

# Statistical mechanical models for quantum codes subject to correlated noise

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# Quantum error correction

Two important questions about quantum codes:

- How do I decode?
- What is the threshold?

The statistical mechanical mapping<sup>1</sup> allows us to address both questions codes subject to Pauli noise.

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# Statistical mechanical mapping

Stabiliser code  
& Pauli noise



Disordered statistical  
mechanical model

Error-correcting  
threshold



Quenched phase  
transition

Maximum likelihood  
decoding



Calculating partition  
functions/free energies

Allows us to reappropriate techniques for studying stat. mech. systems to study quantum codes:

Threshold  
approximation



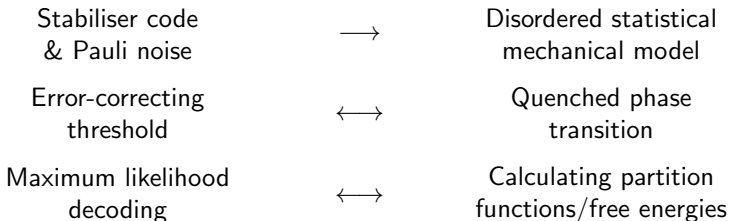
Monte Carlo simulation

Optimal decoding



Partition function  
calculation

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Stabiliser code & Pauli noise	→	Disordered statistical mechanical model
Error-correcting threshold	↔	Quenched phase transition
Maximum likelihood decoding	↔	Calculating partition functions/free energies

Allows us to reappropriate techniques for studying stat. mech. systems to study quantum codes:

Threshold approximation	←	Monte Carlo simulation
Optimal decoding	←	Partition function calculation



# Our results

- Generalise mapping correlated noise for arbitrary codes
- Numerically show that mild correlations can lower the threshold of the toric code considerably
- Show how to apply our mapping to circuit noise via the history code, allowing us to approximate fault tolerant thresholds
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# Independent case

Let  $\llbracket A, B \rrbracket$  be the scalar commutator of two Paulis  $AB =: \llbracket A, B \rrbracket BA$ , and restrict to qubit stabiliser codes.

For a stabiliser code generated by  $\{S_k\}_k$ , and an error Pauli  $E$ , the (disordered) Hamiltonian  $H_E$  is defined

$$H_E(\vec{s}) := - \sum_i \sum_{\sigma \in \mathcal{P}_i} \overbrace{J_i(\sigma)}^{\text{Coupling}} \overbrace{\llbracket \sigma, E \rrbracket}^{\text{Disorder}} \overbrace{\prod_{k: \llbracket \sigma, S_k \rrbracket = -1} s_k}^{\text{DoF}}$$

for  $s_k = \pm 1$ , and coupling strengths  $J_i(\sigma) \in \mathbb{R}$ .

Take-aways:

- Ising-type, with interactions corresponding to single-site Paulis  $\sigma$
- Disorder  $E$  flips some interactions (Ferro  $\leftrightarrow$  Anti-ferro)
- Local code  $\implies$  local stat. mech. model
- Stat. mech. model has a symmetry:  $s_k \rightarrow -s_k$  and  $E \rightarrow ES_k$

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Suppose we have an independent error model

$$\Pr(E) = \prod_i p_i(E_i).$$

Nishimori condition: 
$$\beta J_i(\sigma) = \frac{1}{4} \sum_{\tau \in \mathcal{P}} \log p_i(\tau) \llbracket \sigma, \tau \rrbracket$$

Using the Fourier-like orthogonality relation  $\frac{1}{4} \sum_{\sigma} \llbracket \sigma, \tau \rrbracket = \delta_{\tau, I}$  we get that

$$e^{-\beta H_E(\vec{1})} = \Pr(E).$$

Together with the previous symmetry,

$$Z_E = \sum_{\vec{s}} e^{-\beta H_E(\vec{s})} = \sum_S e^{-\beta H_{ES}(\vec{1})} = \sum_S \Pr(ES) = \Pr(\bar{E}).$$

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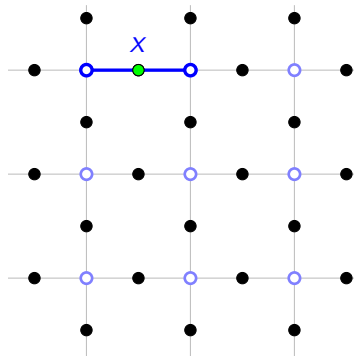
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# Independent examples

## Toric code

Bit-flip  $\rightarrow$  Random-bond Ising<sup>1</sup>

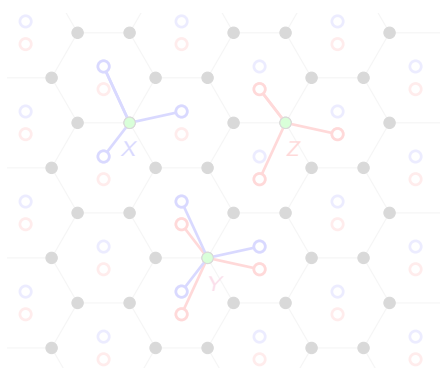
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## Colour code

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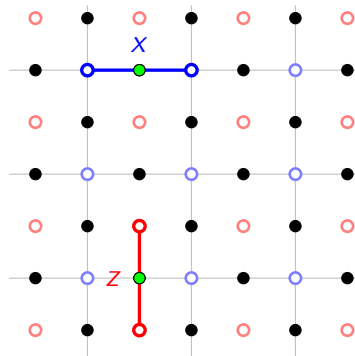
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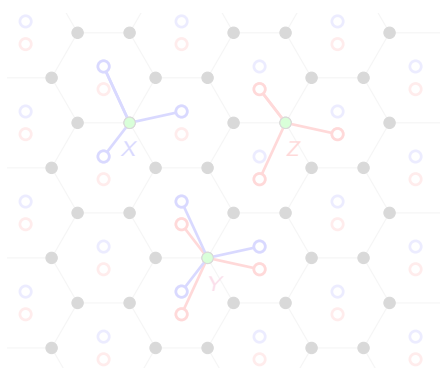
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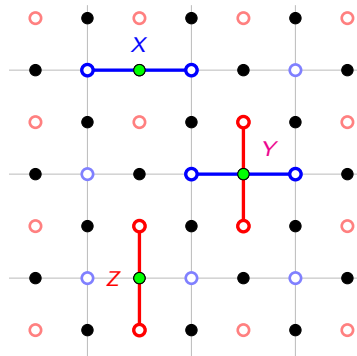
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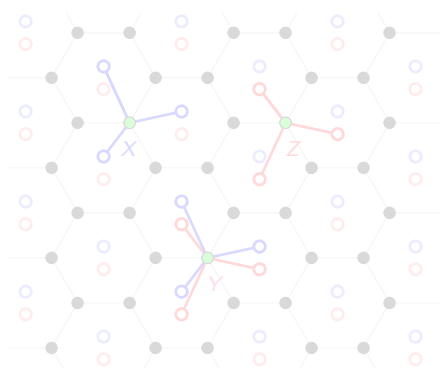
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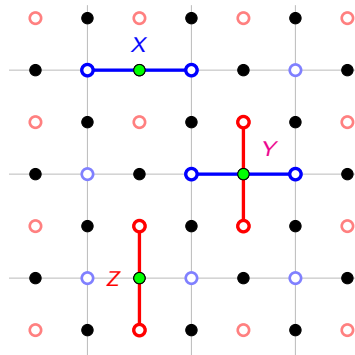
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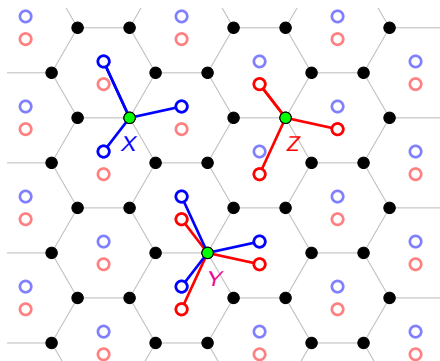
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# Error correction threshold as a quenched phase transition

Consider the free energy cost of a logical error  $L$ ,

$$\Delta_E(L) = -\frac{1}{\beta} \log Z_{EL} + \frac{1}{\beta} \log Z_E.$$

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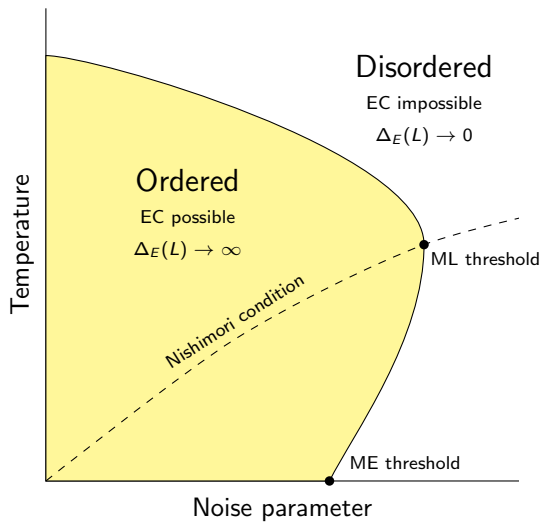
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# Phase diagram sketch





# Correlated case

The key point independence gave us was the ability to factor our noise model

$$\Pr(E) = \prod_i p_i(E_i).$$

We can generalise this to correlated models:

## Factored distribution

An error model factors over regions  $\{R_j\}_j$  if there exist  $\phi_j : \mathcal{P}_{R_j} \rightarrow \mathbb{R}$  such that

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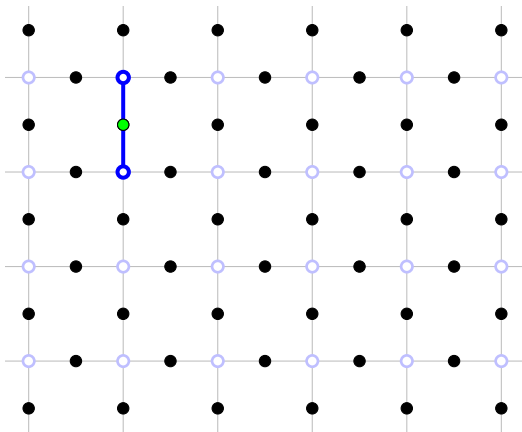
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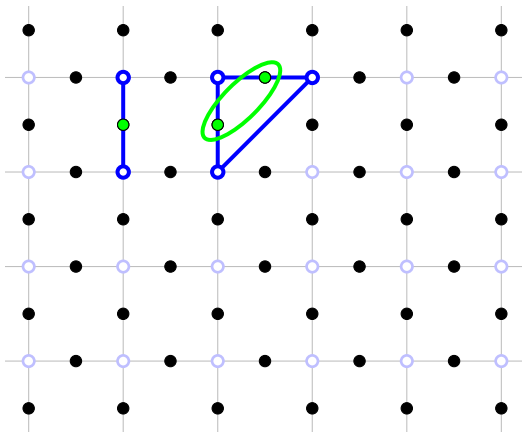
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**Toric code with correlated bit-flips**  
Correlations induce longer-range interactions



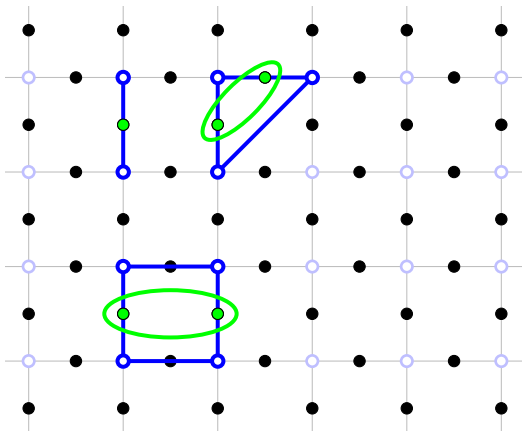
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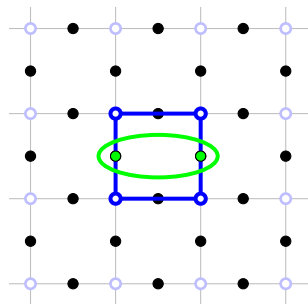
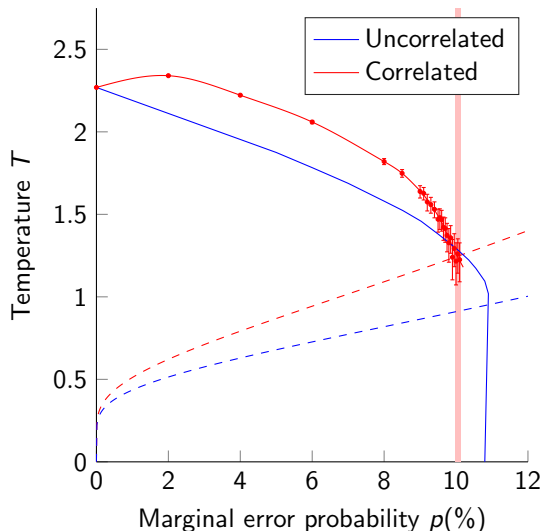
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# Monte Carlo simulations



## Thresholds

Indep.<sup>2,3</sup>:  $p_{th} = 10.917(2)\%$

Corr.:  $p_{th} = 10.04(6)\%$

<sup>3</sup>Dennis et.al., JMP 2002, doi:10/cs2mtf, arXiv:quant-ph/0110143

<sup>3</sup>Toldin et.al., JSP 2009, doi:10/c3r2kc, arXiv:0811.2101

# Decoding

Can the stat. mech. model give us a decoder?

If an error  $E$  occurs, a decoder needs to select one of the logical error classes

$$\overline{E} \quad \overline{EL_1} \quad \overline{EL_2} \quad \overline{EL_3} \quad \dots$$

The optimal (maximum likelihood) decoder selects the most likely class

$$D_{\text{ML}} = \overline{EL_I} \quad \text{where} \quad I = \arg \max_I \Pr(\overline{EL_I}).$$

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# Decoding from partition functions

Along the Nishimori line, the maximum likelihood condition corresponds to maximising the partition function

$$l = \arg \max_l Z_{EL_l}.$$

Approximating  $Z_{EL_l}$  therefore allows us to approximate the ML decoder.

- Step 1: Measure the syndrome  $s$
- Step 2: Construct an error  $C_s$  which has syndrome  $s$
- Step 3: Approximate  $Z_{C_s L_l} = \Pr(\overline{C_s L_l})$  for each logical  $l$
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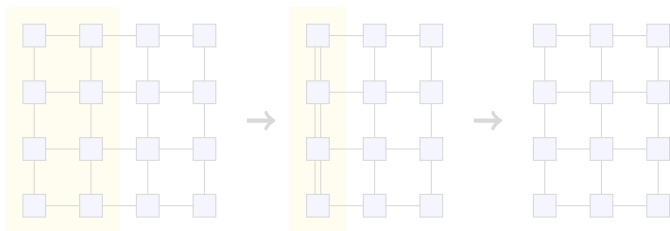
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Partition functions can be expressed as tensor networks<sup>4,5</sup>, allowing us to use approximate tensor network contraction schemes.

For 2D codes and locally correlated noise, this tensor network is also 2D. Here we can use the MPS-MPO approximation contraction scheme considered by Bravyi, Suchara and Vargo<sup>6</sup>:



<sup>4</sup>Verstraete et. al., PRL 2006, doi:10/dfgcz8, arXiv:quant-ph/0601075

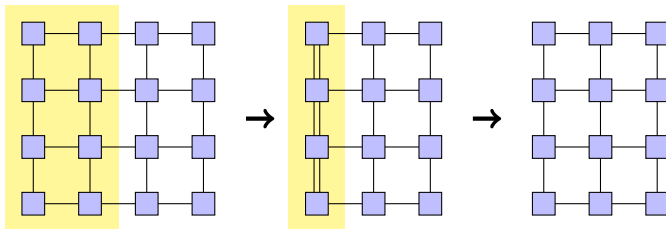
<sup>5</sup>Bridgeman and Chubb, JPA 2017, doi:10/cv7m, arXiv:1603.03039

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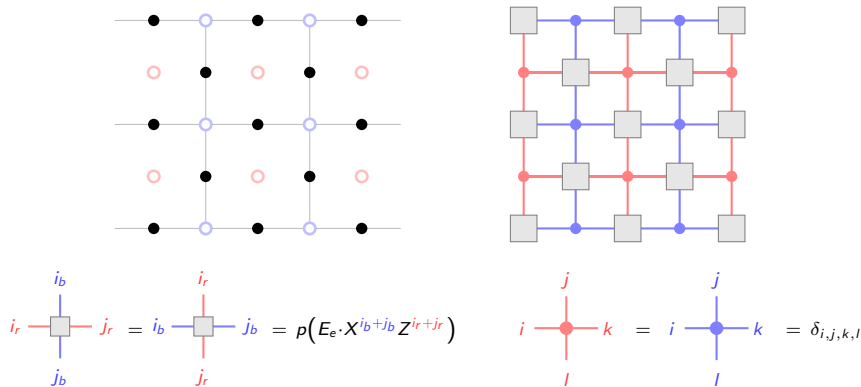
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# Decoding from (approximate) tensor network contraction

This gives an algorithm for (approximate) maximum likelihood decoding for any 2D code, subject to any locally correlated noise, generalising BSV.

Indeed, applying this to iid noise in the surface code reproduces BSV:



# Conclusions and further work

- Extended the stat. mech. mapping to correlated models
  - Can apply stat. mech. mapping to circuit noise via the history code
  - Numerically evaluated the threshold of correlated bit-flips in the toric code
  - Stat. mech. mapping gives tensor network maximum likelihood decoders
- 
- Can we apply this to experimentally relevant correlated models?
  - Can we use the decoders to understand to better understand the connection between correlation and the threshold (ongoing work with David Tuckett and Benjamin Brown).

**Thank you!**

**ArXiv:1809.10704**

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