

Statistical mechanical models and tensor network decoding for quantum codes

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Stat mech mapping (with Steve Flammia, AWS)
arXiv:2101.04125, to appear in AIHP:D

2D tensor network decoding
arXiv:2101.04125

Tensor contraction Julia package
Coming soon!

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SMM+TND

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Quantum codes

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More generally, two important questions arise for any quantum code:

- How do I decode?
- What is the threshold?

The statistical mechanical mapping¹ allows us to systematically address both questions for codes subject to stochastic Pauli noise.

It achieves this by mapping the answers onto the phase boundary of a disordered classical stat mech model.

By combining this with a novel tensor network contraction scheme we get a general and approximately optimal decoder for any 2D code.

¹Dennis, Kitaev, Landahl, Preskill, JMP 2002, doi:10/cs2mtf, arXiv:quant-ph/0110143

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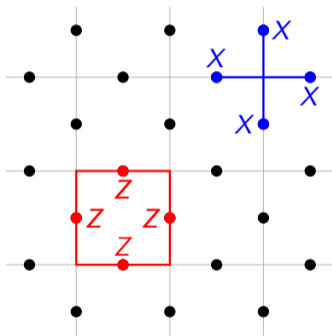
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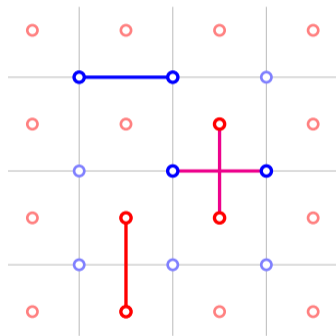
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Toric code



Eight-vertex model

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Stabiliser code
& Pauli noise



Disordered statistical
mechanical model

Threshold



Phase transition

Decoding



Calculating partition
functions

Allows us to reappropriate techniques for studying stat mech systems to study quantum codes, e.g.

Threshold
approximation



Monte Carlo simulation

Optimal decoding



Partition function
calculation

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Stat mech mapping:

- Generalise the mapping to correlated noise for arbitrary subsystem codes
- Numerically study the toric code with mildly correlated errors
- Show how circuit noise can be studied, allowing for fault-tolerant thresholds to be approximated
- Generalise the tensor network decoder of Bravyi, Suchara and Vargo

Tensor network decoding:

- Develop a linearithmic contraction algorithm for 2D tensor networks
- Apply this to several surface/colour codes, including irregular codes
- Provide evidence for a conjecture about surface code thresholds

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Independent case: Hamiltonian

Let $\llbracket A, B \rrbracket$ be the scalar commutator of two Paulis, such that $AB =: \llbracket A, B \rrbracket BA$.

For a stabiliser code generated by $\{S_k\}_k$, and an error Pauli E , the (disordered) Hamiltonian H_E is defined

$$H_E(\vec{s}) := - \sum_i \sum_{\sigma \in \mathcal{P}_i} \overbrace{J_i(\sigma)}^{\text{Coupling}} \overbrace{\llbracket \sigma, E \rrbracket}^{\text{Disorder}} \overbrace{\prod_{k: \llbracket \sigma, S_k \rrbracket = -1} s_k}^{\text{DoF}}$$

for $s_k = \pm 1$, and coupling strengths $J_i(\sigma) \in \mathbb{R}$.

Take-aways:

- Ising-type, with interactions corresponding to single-site Paulis σ
- Disorder E flips some interactions (Ferro \leftrightarrow Anti-ferro)
- Local code \implies local stat mech model

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Independent case: Gauge symmetry

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$$H_E(\vec{s}) = - \sum_i \sum_{\sigma \in \mathcal{P}_i} J_i(\sigma) \llbracket \sigma, E \rrbracket \prod_{k: \llbracket \sigma, S_k \rrbracket = -1} s_k$$

Using $\llbracket A, B \rrbracket \llbracket A, C \rrbracket = \llbracket A, BC \rrbracket$, we see this system has a gauge symmetry

$$s_k \rightarrow -s_k \quad \text{and} \quad E \rightarrow ES_k.$$

This gauge symmetry will capture the logical equivalence of errors, $Z_E = Z_{ES_k}$.

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Independent case: Nishimori condition

Suppose we have an independent error model

$$\Pr(E) = \prod_i p_i(E_i),$$

we now want $Z_E = \Pr(\bar{E})$, \bar{E} is the error class

$$\bar{E} := \{ES | S \in \mathcal{S}\}.$$

Using the gauge symmetry we have that the partition function can be written as a sum stabiliser-equivalent errors

$$Z_E = \sum_{\vec{s}} e^{-\beta H_E(\vec{s})} = \sum_S e^{-\beta H_{ES}(\vec{1})} = \sum_{F \in \bar{E}} e^{-\beta H_F(\vec{1})}.$$

If we select the coupling strength such that $e^{-\beta H_E(\vec{1})} = \Pr(E)$, then $Z_E = \Pr(\bar{E})$ will follow.

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We now want to pick our couplings such that $e^{-\beta H_E(\vec{1})} = \Pr(E)$. Expanding this out, we get

$$\sum_i \log p_i(E) = - \sum_i \sum_{\sigma} \beta J_i(\sigma) \llbracket \sigma, E \rrbracket.$$

Using the Fourier-like orthogonality relation $\frac{1}{4} \sum_{\sigma} \llbracket \sigma, \tau \rrbracket = \delta_{\tau, I}$, this becomes

$$\text{Nishimori condition: } \beta J_i(\sigma) = \frac{1}{4} \sum_{\tau \in \mathcal{P}} \log p_i(\tau) \llbracket \sigma, \tau \rrbracket,$$

which implies $e^{-\beta H_E(\vec{1})} = \Pr(E)$, and therefore $Z_E = \Pr(\bar{E})$.

This intrinsically links the error correcting behavior of the code to the thermodynamic behavior of the model (along the Nishimori line).

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Toric code with iid bit-flips

Step 1: Degrees of freedom

$s_v = \pm 1$ on each vertex v

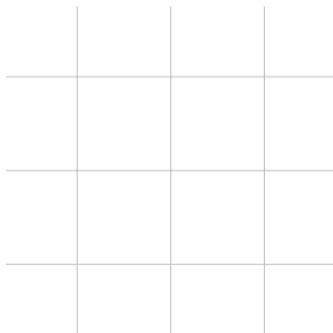
Step 2: Interactions

$$H_I = - \sum_{v \sim v'} J s_v s_{v'}$$

Step 3: Disorder

$$H_E = - \sum_{v \sim v'} J e_{w'} s_v s_{v'}$$

$$\text{where } e_{w'} = \begin{cases} +1 & E_{w'} = I, \\ -1 & E_{w'} = X. \end{cases}$$



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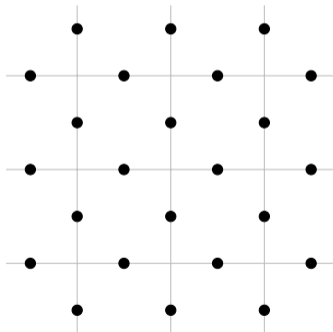
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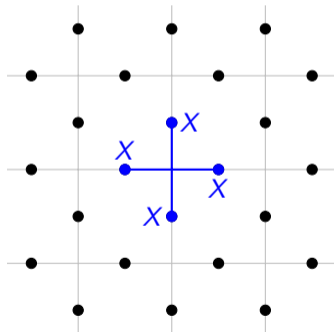
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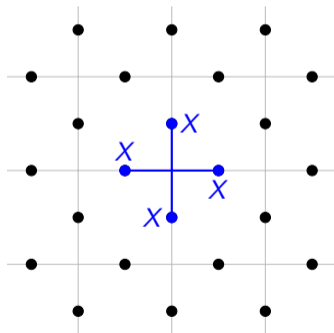
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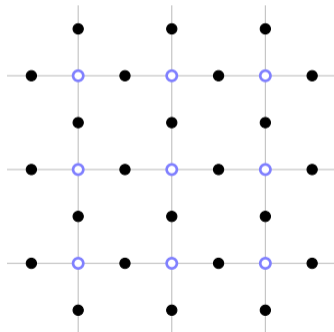
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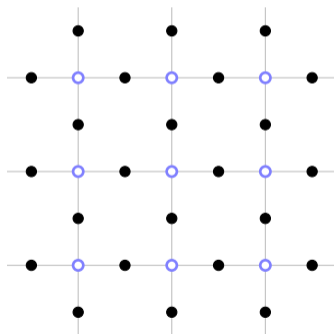
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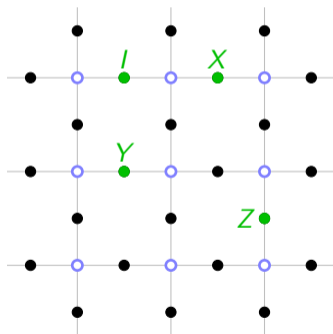
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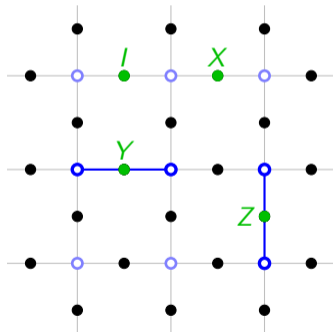
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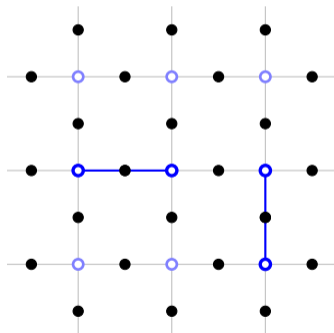
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$$H_I = - \sum_{v \sim v'} J s_v s_{v'}$$

Step 3: Disorder

$$H_E = - \sum_{v \sim v'} J e_{w'} s_v s_{v'}$$

$$\text{where } e_{w'} = \begin{cases} +1 & E_{w'} = I, \\ -1 & E_{w'} = X. \end{cases}$$



Toric code and the random-bond Ising model

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Toric code with iid bit-flips

$$\Pr(X_e) = p, \quad \Pr(I_e) = 1 - p.$$

Step 1: Degrees of freedom

$$s_v = \pm 1 \text{ on each vertex } v$$

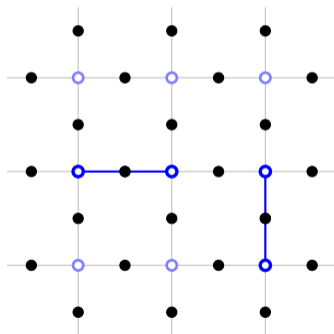
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$$[[\sigma, S_k]] \rightarrow [[\sigma, ES_k]]$$

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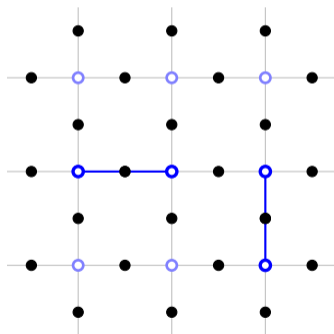
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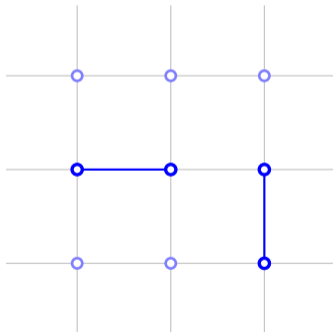
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$$\Pr(+J) = p, \quad \Pr(-J) = 1 - p.$$



$\pm J$ Random-bond Ising Model

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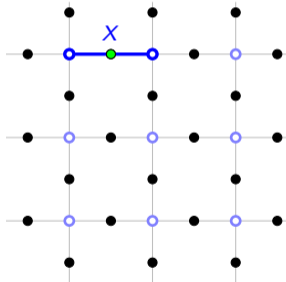
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Toric code

Bit-flip \rightarrow Random-bond Ising¹
Indep. $X&Z \rightarrow 2 \times$ Random-bond Ising
Depolarising \rightarrow Random 8-vertex model²



Colour code

Bit-flip \rightarrow Random 3-spin Ising
Indep. $X&Z \rightarrow 2 \times$ Random 3-spin Ising
Depolarising \rightarrow Random interacting 8-vertex²



¹Dennis et.al., JMP 2002, doi:10/cs2mtf, arXiv:quant-ph/0110143

²Bombin et.al., PRX 2012, doi:10/crz5, arXiv:1202.1852

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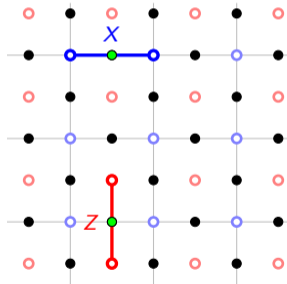
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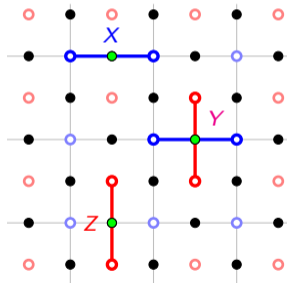
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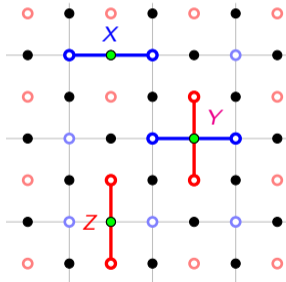
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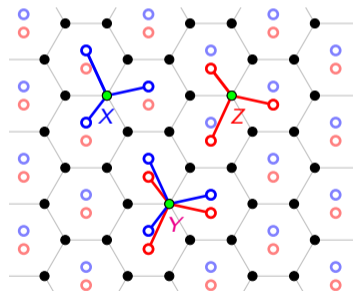
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Error correction threshold as a quenched phase transition

Consider the free energy cost of a logical error L ,

$$\Delta_E(L) = -\frac{1}{\beta} \log Z_{EL} + \frac{1}{\beta} \log Z_E.$$

Along the Nishimori line

$$\Delta_E(L) = \frac{1}{\beta} \log \frac{\Pr(\bar{E})}{\Pr(\bar{EL})},$$

which implies

Below threshold : $\Delta_E(L) \rightarrow \infty$ (in mean)

Above threshold : $\Delta_E(L) \rightarrow 0$ (in prob.)

Error correction threshold as a quenched phase transition

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Phase diagram sketch

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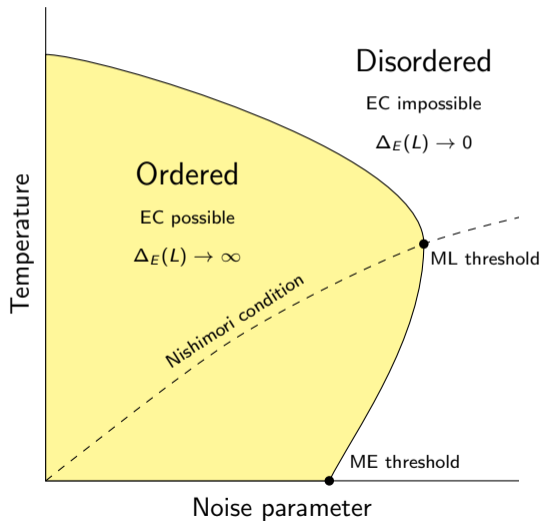
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The key point independence gave us was the ability to factor our noise model

$$\Pr(E) = \prod_i p_i(E_i).$$

We can generalise this to correlated models:

Factored distribution

An error model factors over regions $\{R_j\}_j$ if there exist $\phi_j : \mathcal{P}_{R_j} \rightarrow \mathbb{R}$ such that

$$\Pr(E) = \prod_j \phi_j(E_{R_j})$$

This model includes many probabilistic graphical models, such as Bayesian Networks and Markov/Gibbs Random Fields.

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By construction, we can extend to the correlated case by changing $\sigma \in \mathcal{P}_i$ to $\sigma \in \mathcal{P}_{R_j}$:

$$H_E(\vec{s}) := - \sum_j \sum_{\sigma \in \mathcal{P}_{R_j}} J_j(\sigma) \llbracket \sigma, E \rrbracket \prod_{k: \llbracket \sigma, S_k \rrbracket = -1} s_k$$

Nishimori condition:
$$\beta J_j(\sigma) = \frac{1}{|\mathcal{P}_{R_j}|} \sum_{\tau \in \mathcal{P}_{R_j}} \log \phi_j(\tau) \llbracket \sigma, \tau \rrbracket,$$

As before we get that $Z_E = \Pr(\bar{E})$, and so the threshold manifests as a phase transition.

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Correlated example

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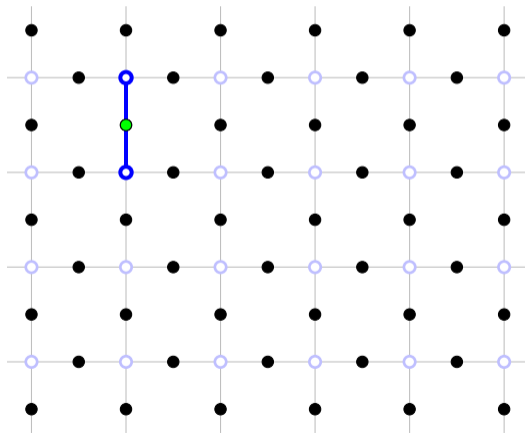
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Toric code with correlated bit-flips
Correlations induce longer-range interactions



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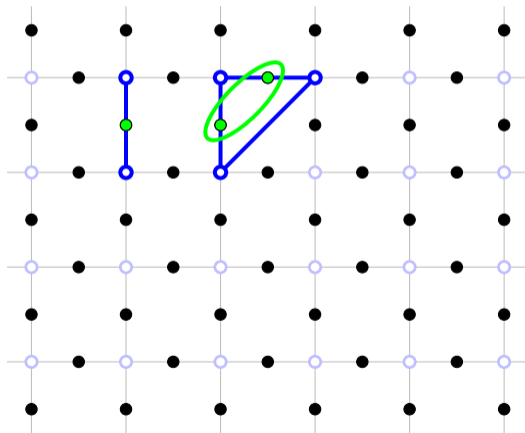
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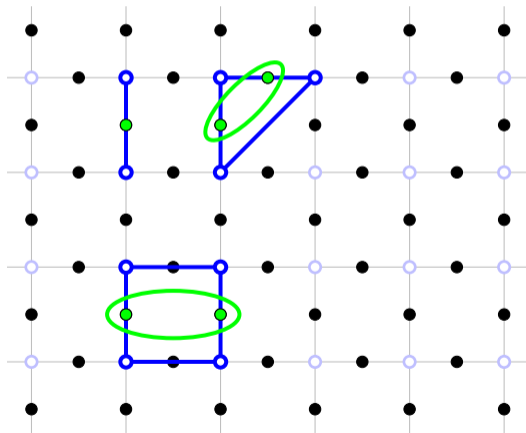
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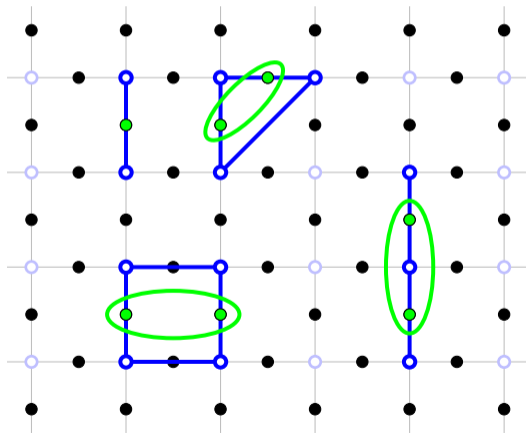
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'Across plaquette' correlated bit-flips

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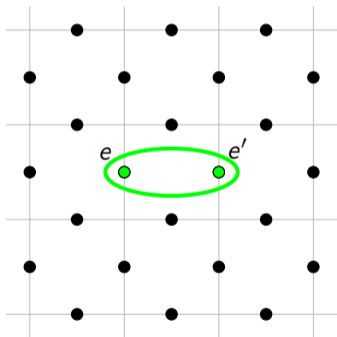
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Model specified by

$$\Pr(I_e | I_{e'}) \quad \Pr(I_e | X_{e'})$$

$$\Pr(X_e | I_{e'}) \quad \Pr(X_e | X_{e'})$$

for all neighbouring edges e and e' .

Convenient parameterisation:

$$\rho := \Pr(X_e), \quad \eta := \frac{\Pr(X_e | X_{e'})}{\Pr(X_e | I_{e'})}.$$

'Across plaquette' correlated bit-flips

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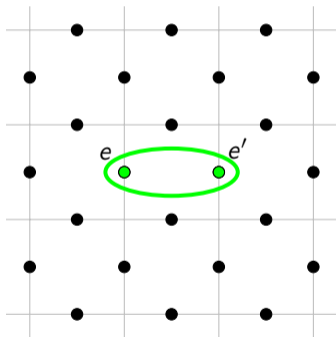
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for all neighbouring edges e and e' .

Convenient parameterisation:

$$p := \Pr(X_e), \quad \eta := \frac{\Pr(X_e|X_{e'})}{\Pr(X_e|I_{e'})}.$$

Monte Carlo simulations

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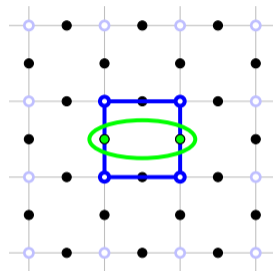
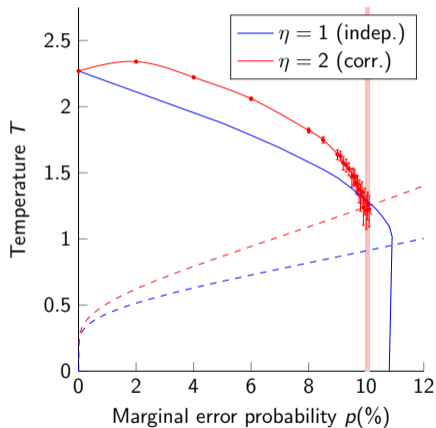
Correlated case

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Thresholds

Indep.: $p_{\text{th}} = 10.917(3)\%^{1,2}$

Corr.: $p_{\text{th}} = 10.04(6)\%$

¹Dennis et.al., JMP 2002, doi:10/cs2mtf, arXiv:quant-ph/0110143

²Toldin et.al., JSP 2009, doi:10/c3r2kc, arXiv:0811.2101

Decoding

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Can the stat mech model give us a decoder?

If an error E occurs, a decoder needs to select one of the degenerate logical error classes

$$\bar{E} \quad \overline{EL_1} \quad \overline{EL_2} \quad \overline{EL_3} \quad \dots$$

The optimal (maximum likelihood) decoder selects the most likely class

$$D_{\text{ML}} = \overline{EL_I} \quad \text{where} \quad I = \arg \max_l \Pr(\overline{EL_l}).$$

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Decoding from partition functions

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Along the Nishimori line, the maximum likelihood condition corresponds to maximising the partition function

$$l = \arg \max_l Z_{EL_l}.$$

Approximating Z_{EL_l} therefore allows us to approximate the ML decoder.

- Step 1: Measure the syndrome s
- Step 2: Construct an arbitrary error C_s which has syndrome s
- Step 3: Approximate $Z_{C_s L_l} = \Pr(\overline{C_s L_l})$ for each logical l
- Step 4: Find the l such that $Z_{C_s L_l}$ is maximised
- Step 5: Apply $(C_s L_l)^{-1}$

Decoding from partition functions

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ML decoding as a tensor network

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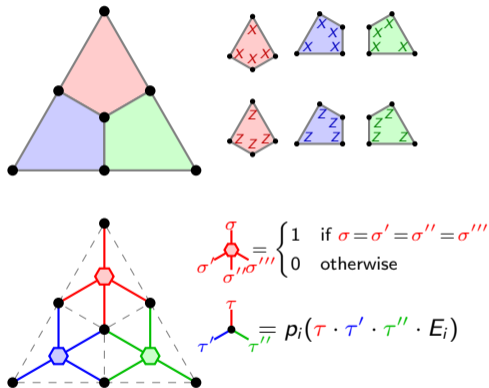
Correlated case

Tensor network decoding

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- Partition functions, and thus error class probabilities, can be expressed as tensor networks^{2,3}.
- There is a tensor for each qubit, and for each stabiliser
- By approximating these class probabilities we can find the ML error

²Verstraete et. al., PRL 2006, doi:10/dfgcz8, arXiv:quant-ph/0601075

³Bridgeman and Chubb, JPA 2017, doi:10/cv7m, arXiv:1603.03039

Sweeping contraction algorithm

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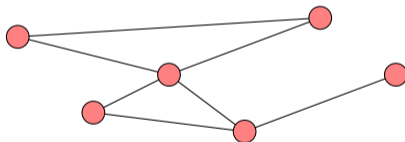
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- We sweep along the network, using lossy compression of Matrix Product States
- To do this for general graphs we need to use a sweeping approach



I'm going to release a general purpose Julia implementation of this soon.

Sweeping contraction algorithm

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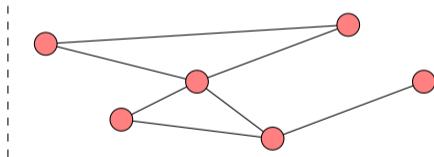
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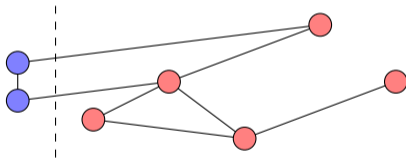
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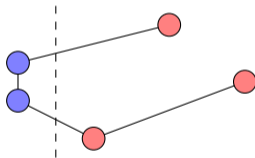
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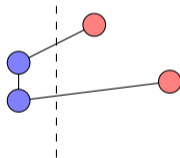
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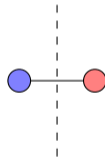
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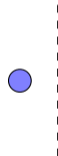
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- We sweep along the network, using lossy compression of Matrix Product States
- To do this for general graphs we need to use a sweeping approach

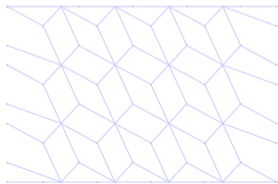
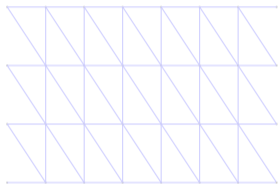


I'm going to release a general purpose Julia implementation of this soon.

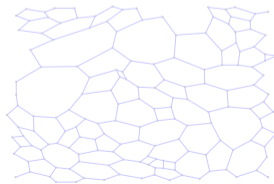
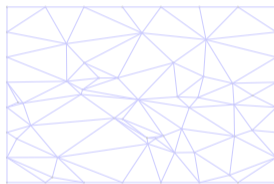
Surface codes on different lattices

How does the surface code perform on different lattices/graphs?

Raise or lower the connectivity



Irregular graphs



SMM+TND

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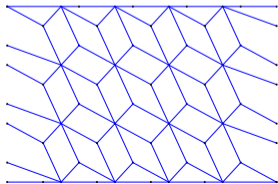
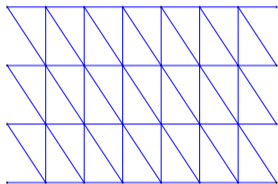
Results

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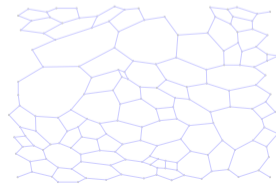
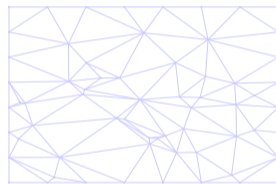
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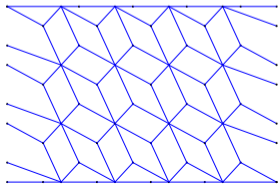
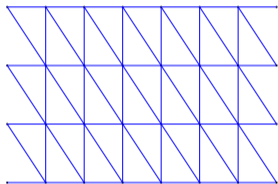
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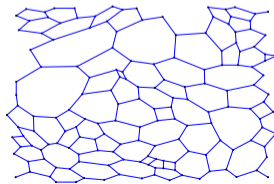
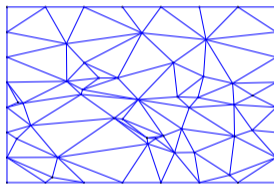
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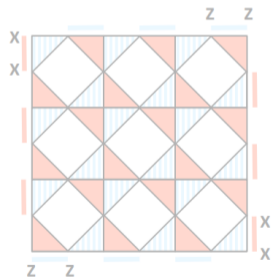
Conclusion

We also want to capture 2D models that aren't (ostensibly) surface codes.

Colour code



Subsystem surface code



Non-SC topological models

SMM+TND

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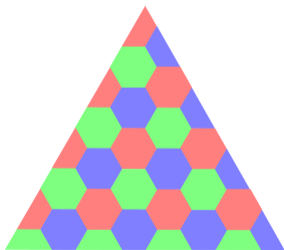
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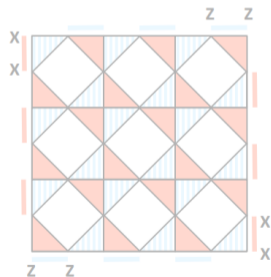
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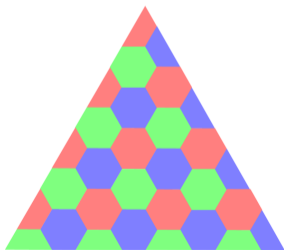
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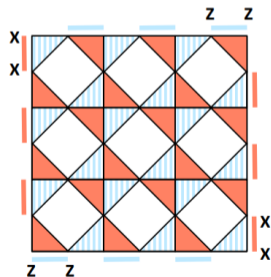
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Thresholds

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	Bit-flip		Phase-flip		Depolarising	
	Observed	Upper bound	Observed	Upper bound	Observed	Upper bound
Surface code (reg.)						
Square	10.917(5)%	10.9187%	10.917(5)%	10.9187%	18.81(3)%	18.9(3)%
Tri./Hex.	16.341(7)%	16.4015%	6.748(5)%	6.7407%	13.81(7)%	?
Kag./Rho.	9.875(5)%	?	11.910(6)%	?	18.09(4)%	?
T.H./Asa.	4.297(7)%	?	20.701(13)%	?	9.07(8)%	?
Surface code (irr.)						
Rand. Tri.	17.128(15)%	?	6.237(9)%	?	12.85(3)%	?
Rand. Quad.	12.195(12)%	?	9.715(11)%	?	18.05(3)%	?
Subsystem SC						
Square	6.705(13)%	6.7407%	6.705(13)%	6.7407%	11.23(3)%	?
Colour Code						
Hexagonal	10.910(5)%	10.9(2)%	10.910(5)%	10.9(2)%	18.68(2)%	18.9(3)%

Surface code saturating the Hashing Bound

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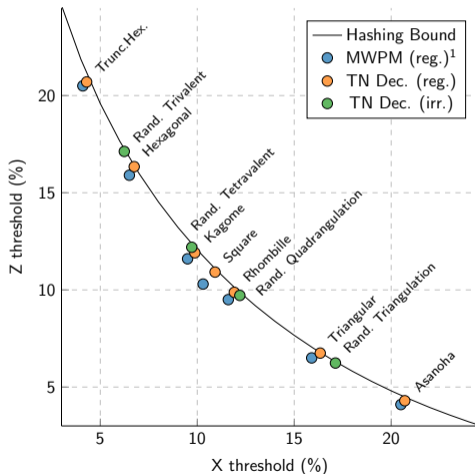
Conclusion

We find a trade-off between the X and Z thresholds.

Thresholds saturate the hashing bound

$$h(\tau_x) + h(\tau_z) \leq 1$$

Pair matching studied earlier by Fujii et.al.⁴



⁴Fujii et.al., doi.org/d5sb, arXiv:1202.2743

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Pair matching studied earlier by Fujii et.al.⁴

	Entropy	
	MWPM ¹	TN
Regular		
Square	0.957	0.9948(3)
Tri./Hex.	0.979	0.9989(3)
Kag./Rho.	0.971	0.9918(3)
T.H./Asa.	0.979	0.9915(6)
Irregular		
Rand. Tri.	?	0.9974(7)
Rand. Quad.	?	0.9948(7)

⁴Fujii et.al., doi.org/d5sb, arXiv:1202.2743

Future work

- Stat mech mapping beyond Pauli codes
- Using the stat mech mapping for more than the threshold
- TN decoding of correlated noise
- TN decoding in 3D (including noisy 2D) and beyond, LDPCs, etc.

Thank you!

Stat mech mapping: arXiv:1809.10704, to appear in AIHPD

Tensor network decoding: arXiv:2101.04125

Tensor contracting package: Coming soon!

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🔗 christopherchubb.com

🐦 @QuantumChubb

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